

# Rayleigh scattering from argon clusters in a planar expansion

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#### Abstract

Rayleigh scattering is presented as the evidence for the presence of large argon clusters formed in a planar expansion. Based on the observed scattering signal, the dependence of mean cluster size on stagnation pressure is  $\langle N \rangle \propto P_0^{3.38}$ . This is in contrast to the dependence of the mean cluster size on stagnation pressure for a symmetric expansion of  $\langle N \rangle \propto P_0^{2.29}$ . Using interferometry in conjunction with the Rayleigh scattering signal we are able to estimate the mean cluster size for clusters formed in the planar expansion.

#### Principle of Corresponding Jets

The <u>principle of corresponding jets</u> is used to predict mean cluster size given a variety of parameters:

$$\Gamma^* = \frac{kd^{*q}}{T_0^{2.5-sq}} P_0$$

 $\Gamma^*$  is a dimensionless parameter related to the cluster size

k is a characteristic constant that depends upon the gas the geometry of the expansion.

 $d^*$  is the effective aperture opening

q varies from 0 to 1 and is determined experimentally by varying d and examining cluster size. It is 0.85 for axisymmetric expansions.

s depends upon geometry and is 0.25 for axisymmetric expansions and 1 for a planar expansion.

The mean cluster size is determined by  $\langle N \rangle = C \Gamma^{*\beta}$ 

## Measuring the mean cluster size by Rayleigh scattering

Differential Rayleigh scattering cross section:

$$\frac{d\sigma}{d\Omega}\Big|_{total} = \frac{9\pi^2 M_a^2}{\rho^2 \lambda^4} \left(\frac{n^2 - 1}{n^2 + 2}\right)^2 \left[N_c \left\langle N^2 \right\rangle\right] \sin^2 \phi$$

Vol<sup>2</sup> =  $\frac{M_a^2}{\rho^2} = \frac{16\pi^2}{9} R^6$ ,  $M_a$  is the atomic mass,  $\rho$  is the density of

the cluster, *R* is the radius of a cluster.

*n* is the index of refraction of the cluster. ( $\sim$ 1.3 for argon clusters)

 $N_c$  is the number density of clusters

 $\langle N^2 \rangle$  is the square of the mean cluster size.

 $\phi$  is the angle between the incident radiation's polarization and the polarizability of the cluster.

The number density of clusters is unknown but can be related to the

tomic number density by: 
$$N_c = \left(\frac{\alpha N_a}{\langle N \rangle}\right)$$

 $N_a$  is the atomic number density in the expansion and is proportional to the stagnation pressure  $P_o$ .

 $\alpha$  is the fraction of atoms in clusters.

$$\langle N \rangle$$
 is the mean cluster size  $\frac{\langle N^2 \rangle}{\langle N \rangle} \sim \langle N \rangle$  and  $\langle N \rangle = C \Gamma^* {}^{\beta} \propto P_0{}^{\beta}$ 

 $\sigma |_{total} \propto P_0 \langle N \rangle \propto P_0^{\beta+1}$ 

By determining the dependence of Rayleigh scattering signal on stagnation pressure we can find the dependence of the mean cluster size on stagnation pressure determining  $\beta$ .



before the detector, PMT is the photomultiplier tube, BD is a beam dump, and P is the pressure sensor inside the plenum.



Pressure dependence of Rayleigh scattering for a planar expansion



The peak value of the scattered signal for the planar expansion is plotted as a function of measured stagnation pressure. The uncertainty in the values is no larger than the size of the markers. The fit curve correspondZing to  $S_{R} = Dp_{0}^{4.38\pm0.09}$  ( $\beta = 3.38\pm0.09$ ) is shown.

Pressure dependence of Rayleigh scattering for an axisymmetric expansion



The peak value of the scattered signal for the axisymmetric expansion is plotted as a function of measured stagnation pressure. The uncertainty in the values is no larger than the size of the markers. The fit curve corresponding to  $S_{R} = Dp_{0}^{3.29 \pm 0.06}$  ( $\beta = 2.29 \pm 0.06$ ) is shown. Reported values range from  $\beta$ =1.5 to 2.5



Relating this to the cluster size we get  $\vec{P} =$ 

Given a size distribution of clusters  $N_c(N_i)$ , we can change the summation to be over cluster size rather than molecule so that

$$\vec{P} = \left[ \sum_{\substack{\text{cluster}\\\text{size}}} \left( \frac{\varepsilon - 1}{\varepsilon + 2} \right) \left( \frac{3}{4\pi} \frac{M_a}{\rho} N_c N_i N_i \right) \right] \vec{E}$$

Completing the summation results in  $\vec{P} = \left[ \left( \frac{\varepsilon - 1}{\varepsilon + 2} \right) \frac{3}{4\pi} \left\{ \frac{M_a}{\rho} \langle N \rangle N_c \right\} \right] \vec{E}$ 



In the Rayleigh scattering cross section, we did not know the number of clusters. Comparing this to the total detected Rayleigh scattered light:



gas we find:

To determine the cluster size in this way you must: •Have an absolute calibration of detection system. •Accurately measure the phase shift due to the gas.

## Interferometry Experimental Configuration



## Interferometry

The polarization of dielectric sphere in an external field is given by:

$$\sum_{\substack{i \\ molec}} \left(\frac{\varepsilon - 1}{\varepsilon + 2}\right) \left(\frac{3}{4\pi} \frac{M_a}{\rho} N_i\right) \vec{E}$$

Since  $\vec{D} = \vec{E} + 4\pi \vec{P} = \varepsilon \vec{E}$ , the relative dielectric constant

 $\varepsilon = n^2 = \left( 1 + 4\pi \left[ \left( \frac{\varepsilon - 1}{\varepsilon + 2} \right) \frac{3}{4\pi} \left\{ \frac{M_a}{\rho} \langle N \rangle N_c \right\} \right] \right) \text{ yielding an index of refraction}$ 

Note that the change in index of refraction is proportional to the product of  $\langle N \rangle N_c$  the product of the mean cluster size and the number of clusters.

$$\left[\int \frac{9\pi^2 M_a^2}{\rho^2 \lambda^4} \left(\frac{n_b^2 - 1}{n_b^2 + 2}\right)^2 \left[N_c \left\langle N^2 \right\rangle\right] \sin^2 \phi \, d\Omega\right] \frac{P_0}{A}$$

So if we look at the ratio scattering signal to the index of refraction of the

#### Issues

BS is the beam splitter, M are mirrors, PD is a photodiode to detect fringes, G is the gas line, A is an aperture, IF is an interference filter/aperture before the detector, PMT is the photomultiplier tube, BD is a beam dump, and P is the pressure sensor inside the plenum.

